THE APPLICATION OF IMPLICIT TECHNIQUES IN OCEAN MODELING

Erik Bernsen and Henk A. Dijkstra IMAU, Department of Physics and Astronomy, Utrecht University, The Netherlands

Fred Wubs and Jonas Thies Department of Mathematics and Computer Science University of Groningen, The Netherlands

Annual mean surface wind velocity

FERRET (no GUI) Van 4.3D
--- NOAA/PNEL TWAP
- Jan - 5 1997 15:20:17

Surface Ocean Circulation

Wind-driven circulation: circulation associated with direct forcing of the wind

After: Sverdrup, H.U. et al. (1942)

Temperature + Salinity section

Temperature [°C]

Annual mean surface heat flux

FERRET (no GUI) Van 4.3D
--- NOAA/PNEL TWAP
- Jan - 5 1997 15:00:02

NET DOWNWARD HEAT FLUX (W/M/M)

Annual mean freshwater flux

Annual net precipitation - evaporation

Global Conveyor Circulation

Ganachaud & Wunsch, Nature, 408, 453, (2000).

QuickTime™ and a Cinepak decompressor are needed to see this picture.

 $1 Sv = 10^{6} m^{3/s}$

Thermohaline circulation: circulation associated with the transport of heat and salt

Meridional Overturning Circulation (MOC) $\boldsymbol{\mathcal{S}}$ ThermoHaline Circulation (THC)

Meridional overturning streamfunction Model determined

MOC: Total northward/southward transport in latitude/depth (is observable) THC: Part of MOC driven by heat/freshwater exchange at the surface and subsequent vertical mixing (is an interpretation)

QuickTime™ and a Sorenson Video decompressor are needed to see this picture.

Conceptual Model of the Atlantic MOC

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The salt-advection feedback

Multiple equilibria in an Atlantic ocean model

Source: Bryan, F. High-latitude salinity effects and interhemispheric thermohaline circulations, Nature, 323, 301-30

I.M.A.U., Utrecht University 11 and 17/08 **11** and 18 and 19 and 19 and 11

... and in global ocean models and EMICs

Rahmstorf et al., GRL, (2005)

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Use of implicit methods

• I: Equilibria versus parameters ★ Steady state solvers and bifurcation analysis ✦continuation: no transient integration

• II: Spin-up

★Use of Newton-Raphson like techniques in explicit models ✦equilibration time scale ~ 5000 yr time step 1 hr

Ocean model formulation

Local conservation of mass, momentum, heat and salt:

$$
\overline{\mathcal{M}_{\lambda}\frac{\partial \mathbf{u}}{\partial t}+\mathcal{L}_{\lambda}\mathbf{u}+\mathcal{N}_{\lambda}(\mathbf{u})\mathbf{u}}=\mathbf{F}_{\lambda}}
$$

 $\overline{\mathcal{M}_\lambda,\mathcal{L}_\lambda,\mathcal{N}_\lambda}$: Operators

: Parameter Vector

: State Vector

+ Boundary conditions & Initial conditions

Governing equations

$$
\frac{Du}{dt} - uv \tan \theta - 2\Omega v \sin \theta + \frac{1}{\rho_0 r_0 \cos \theta} \frac{\partial p}{\partial \phi} =
$$
\n
$$
A_V \frac{\partial^2 u}{\partial z^2} + A_H L_u(u, v) + \frac{\tau_0}{\rho_0 H_m} \tau^{\phi} \mathcal{G}(z)
$$
\n
$$
\frac{Dv}{dt} + u^2 \tan \theta + 2\Omega u \sin \theta + \frac{1}{\rho_0 r_0} \frac{\partial p}{\partial \theta} =
$$
\n
$$
A_V \frac{\partial^2 v}{\partial z^2} + A_H L_v(u, v) + \frac{\tau_0}{\rho_0 H_m} \tau^{\phi} \mathcal{G}(z)
$$
\n
$$
\frac{\partial p}{\partial z} = \rho_0 g(\alpha_T T - \alpha_S S)
$$
\n
$$
\frac{\partial w}{\partial z} + \frac{1}{r_0 \cos \theta} \left(\frac{\partial u}{\partial \phi} + \frac{\partial (v \cos \theta)}{\partial \theta} \right) = 0
$$
\n
$$
\frac{DT}{dt} - \nabla_H \cdot (K_H \nabla_H T) - \frac{\partial}{\partial z} \left(\frac{K_V}{\partial z} \right) = \frac{(T_S - T)}{\tau_T} \mathcal{G}(z)
$$
\n
$$
\frac{DS}{dt} - \nabla_H \cdot (K_H \nabla_H S) - \frac{\partial}{\partial z} \left(\frac{K_V}{\partial z} \right) = F_0 F_S \mathcal{G}(z)
$$

Continuation Methods

$$
M\frac{dx}{dt} + G(x, \lambda) = 0, x \in R^d
$$

d: # degrees of freedom

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A tailored solver for bifurcation analysis of ocean-climate models

Arie de Niet^a, Fred Wubs^{a,*}, Arjen Terwisscha van Scheltinga^b, Henk A. Dijkstra b

Journal of Computational Physics 227 (2007) 654–679

- Full 3D primitive equations
- Solutions for ocean models with $d \sim 1,250,000$
- Efficient handling of bathymetry
- Implementation of state of the art mixing schemes, such as neutral physics and GM

Recent developments: Parallel implementation into TRILINOS

Bifurcation diagram (global ocean model)

II. Spin-up

Spinup time: The (CPU) time needed for the ocean model to reach an equilibrium.

Spinup timescale (physical time): Set by vertical diffusivity K and mean depth of the ocean basin D

Jacobian Free Newton-Krylov (JFNK) methods

Goal: Find a steady state of:

$$
\mathrm{d}\vec{x}/\mathrm{d}t=F(\vec{x})
$$

Timestepper: Iterate

until $\|\vec{x}_{k+1}-\vec{x}_k\|$ small.

$$
\vec{x}_{k+1} = \vec{x}_k + \Delta t F(\vec{x})
$$

Newton-Raphson: Iterate until $\|F(\vec{x}_k)\|$ small

> $\overline{\mathcal{L}}$ *J*

$$
0=F(\vec{x}_k)+J\,\left(\vec{x}_{k+1}-\vec{x}_k\right)
$$

During GMRES method:

$$
\vec{v} \approx \frac{\vec{F}(\vec{x}_k + \epsilon \vec{v}) - \vec{F}(\vec{x}_k)}{\epsilon}
$$

(with ε small)

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A method to reduce the spin-up time of ocean models

Erik Bernsen^{a,*}, Henk A. Dijkstra^a, Fred W. Wubs^b

Ocean Modelling 20 (2008) 380-392

- 3D planetary geostrophic model; T and S only prognostic variables (Samelson & Vallis, 1997)
- Idealized wind-stress and restoring temperature and salinity forcing
- Preconditioner based on dependencies in Jacobian
- Continuation technique in model forcing to avoid problems in Newton-Raphson convergence

Recent developments: Implementation into ocean GCM (MOM4)

'Normal' Spin-up run (16x16x18, PGOM)

Comparison of methods (16x16x18)

Speedup of JFNK method

Conclusion and future work

Implicit techniques have a great application potential in ocean modeling:

Bifurcation analyses possible for global ocean-climate models with hor^{2°} ntal resolution.

★ Future: Improved tailored parallel preconditioners

JFNK methods can provide efficiency for spin-up problems using explicit models

★ Future: Application to statistical steady states